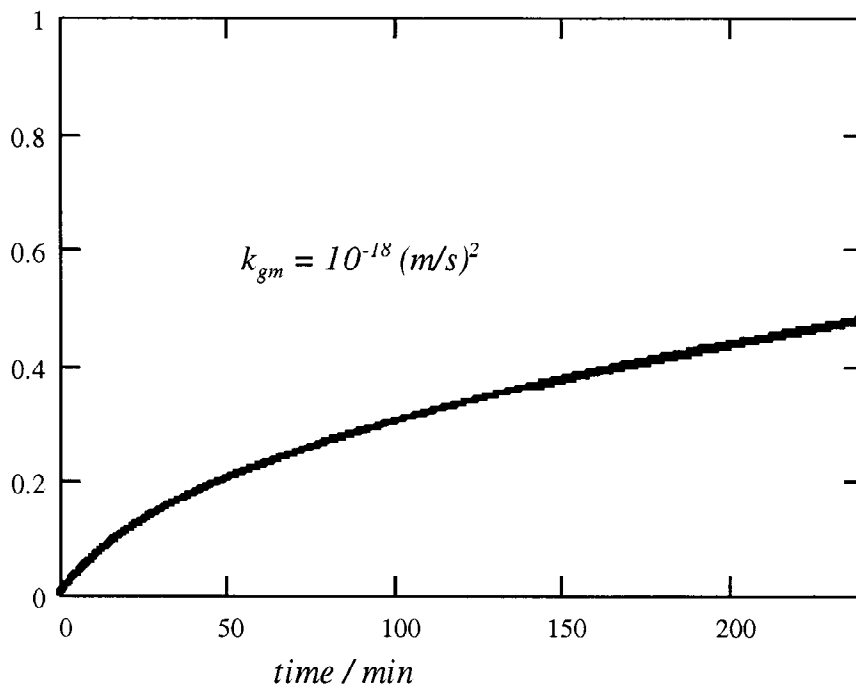


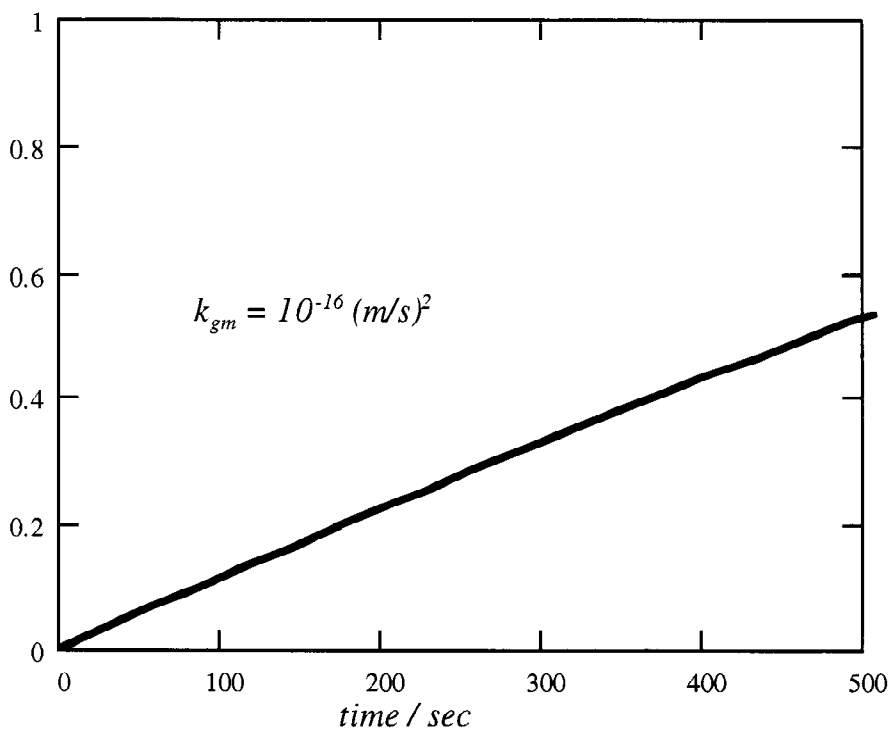
# **EFFECT OF VERY LOW FREQUENCY GROUND MOTION ON THE LHC**

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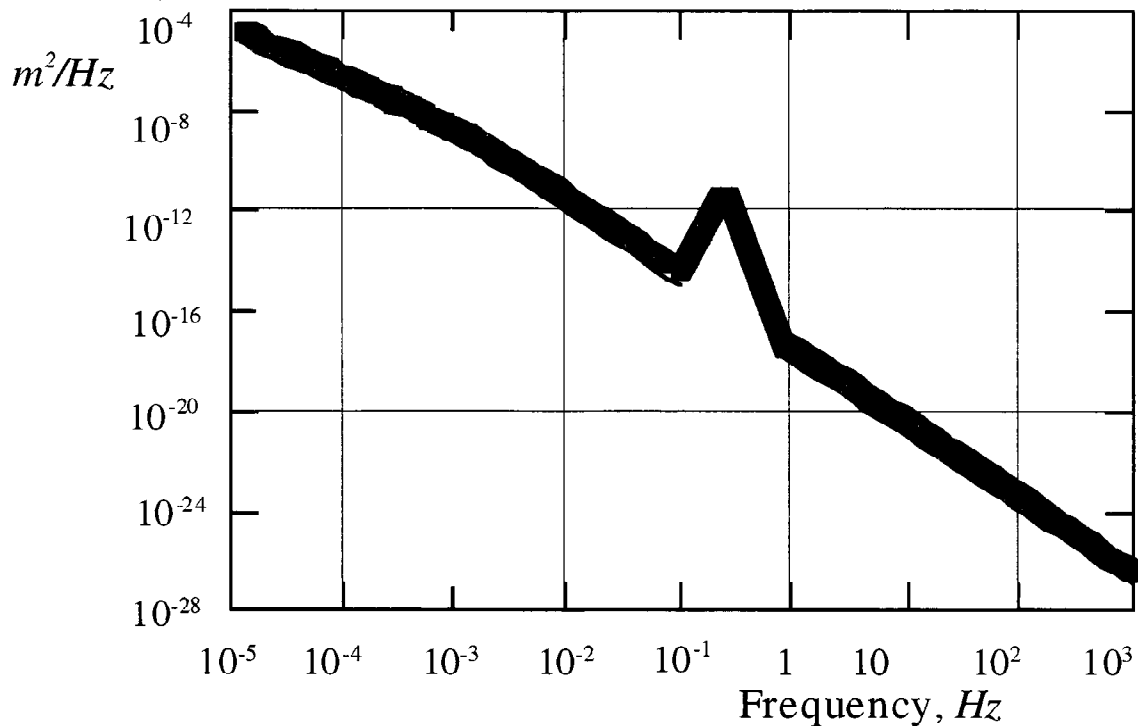
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# 1 INTRODUCTION

Power spectrum



(cultural noise taken out)

- power of ground vibrations increases steeply with decreasing frequency

====> possibility of non-negligible orbit deformations if the motion of accelerator quadrupoles is uncorrelated

- what about **plane wave excitation** ?

ground motion wavelength matches the betatron wavelength

====> very **narrow band** effect for frequencies around 1 Hz, *small spectral power*, small separation effect in LHC  $< 1/1000$  of *rms* beam size in spite of non-negligible optical amplification factor

### Introductory remarks

- non-correlated, or partially correlated *movements*, are observed at short distances in spite of much longer wavelengths
- *orbit changes* related to this can be measured in a large machine as LEP (consequences are harmless, single bore)

but, orbits in the LHC of the two constituent rings may wander apart and hence partially separate the beams

====> *loss of luminosity + modification of beam-beam interaction* if no adequate correcting action is taken

- direct observation of low frequency non-correlated motion led to formulation of the *ATL* scaling law (1991),

**not valid for frequencies above  $\sim 1$  mHz**

====>

a **model** is proposed based on geo-physical arguments valid for the full frequency range

## 2 GROUND MOTION MODEL

- ocean well spectral peak is coherent

the high-pass cut-off frequency of  $\sim 0.2 \text{ Hz}$  together with the speed in water ( $1.5 \text{ kms}^{-1}$ ) suggest a limiting wavelength in the oceans of around  $\sim 7 \text{ km}$   $\sim$  depth of the abyssal plain (between 3 and 5.5 km)

====> these waves are **surface** waves which will **not contribute** to random motion (difficult to imagine how geological fault structures can cause these waves to loose coherence over a fraction of a wavelength)

- remaining spectrum tends to fall with frequency as  $f^{-3}$  above the *ocean hum*

- spectrum falls as  $f^{-2}$  well below this frequency (wavelengths involved exceed 25 km)

- clear evidence exists on lack of correlation (randomness) of very low frequency noises at distances much less than the wavelength

- *ATL* scaling law matches the low frequency  $f^{-2}$  slope and states :

*the random (integrated) motion between two points is proportional to their distance  $L$  while the proportionality factor  $A$  depends on the local properties concerning the randomness*

However, this scaling law yields non-physical results above a certain frequency: it predicts differential movements that are much larger than the absolute ones

====> a different model is proposed that fulfills two requirements

- 1 a general formulation of the motion of a single point, in principle valid everywhere
- 2 take into account local randomisation of the low frequency earth movements.

## 2.2 Basic model

The model consists of a transfer function and a source of excitation

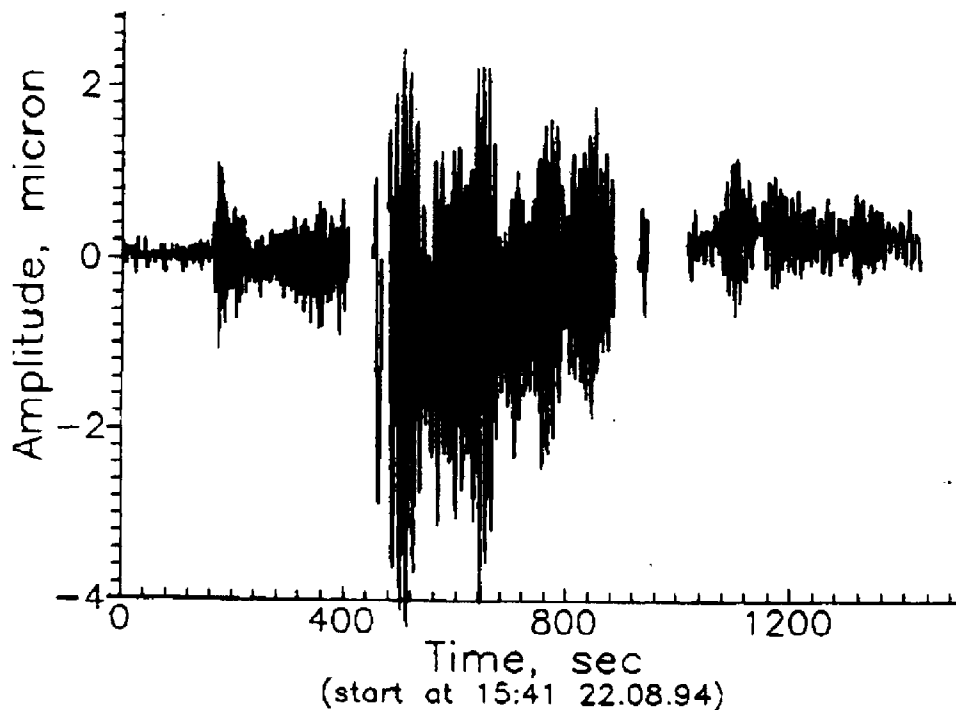
### 2.2.1 transfer function

maximum seismic length of the earth is  $\sim 1500$  s

seismic 'depth' of the earth is about 1/3 of this

====> a cut-off frequency of  $f_{co} \sim 2$  mHz

confirmed by the far away amplitude response of earth quakes:



fast oscillation is the response of the oceans pounding on the continents

amplitude response is not incompatible with a high-pass filter behaviour with a power cut-off frequency  $f_{co} \sim 1.5$  mHz

note : seismic wave attenuation with distance is very small

### 2.2.2 source

the obvious source is : *earth quakes!*

the transfer model suggests a source with a  $f^{-3}$  frequency slope to reproduce noise spectra

examination of a substantial body of phenomenological material has led to the *Gutenberg- Richter law* :

$$\log(n) = -M,$$

$n$  is the number of earth quakes in a given area with a magnitude  $M$  or larger

$$M = \frac{1}{m} \log\left(\frac{x}{x_0}\right)^2$$

$x_0$  is the lower observation limit of seismographs.

According to *Hamblin*  $m = 2$  so that :

$$\left(\frac{x}{x_0}\right)^2 = \frac{1}{n^2}.$$

The power density of the source can be found by differentiation :

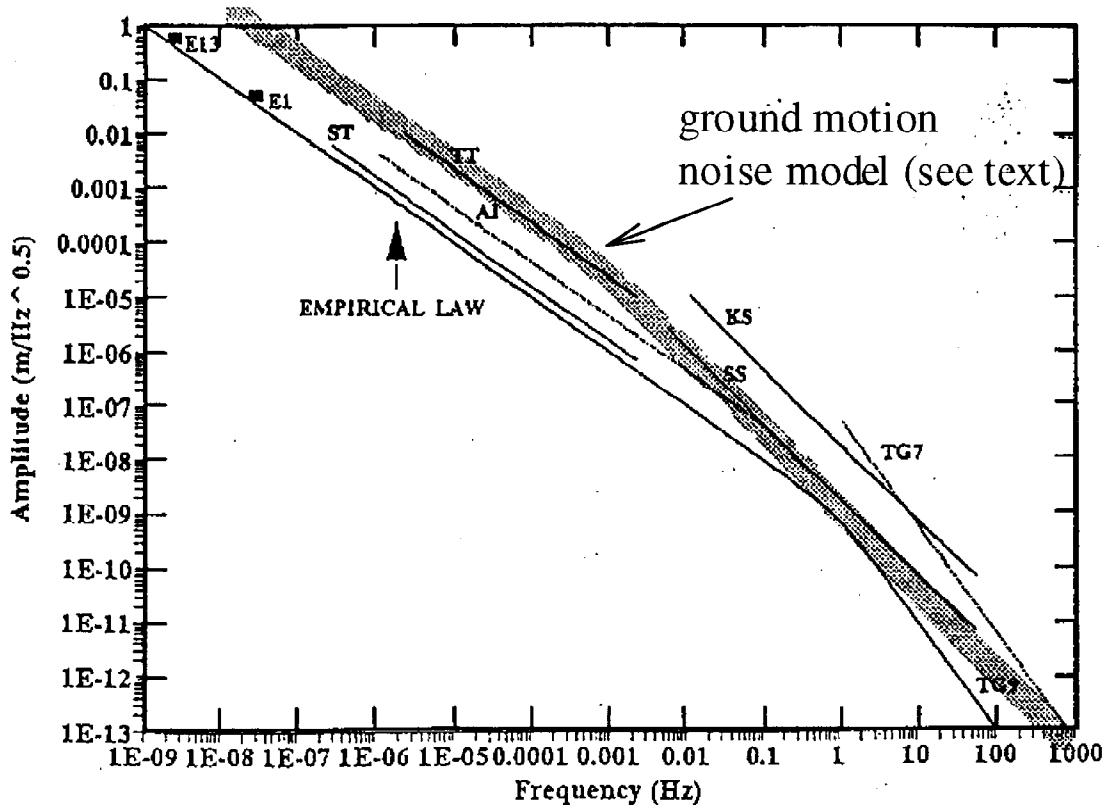
$$\frac{d(x/x_0)^2}{dn} = -\frac{2}{n^3} \propto -\frac{1}{f^3}$$

Combining transfer function and source of vibrations yields the ground motion spectrum :

$$\frac{dx^2}{df}(f) = \frac{k_{gm}}{f^2 \sqrt{f_{co}^2 + f^2}} [m^2/Hz].$$

$k_{gm}$  is a proportionality factor of the source to fit the observed spectral power density. It is a **non-local quantity** that varies from  $\sim 10^{-18} m^2/s^2$  to  $\sim 10^{-16} m^2/s^2$  depending on the state of **global excitation**

**Comparison with various observations :**



The line marked 'empirical law' is related to the model proposed by Takeda e.a..

## 2.2 Randomness

- **basic question :**

how can two points, close together ( $\sim 100\text{ m}$ , much less than a wavelength : many  $\text{km}$ ), move independently?

- **combination of observations may help to explain :**

1 earth quake observations show that depth of the sources is very often  $\sim 30\text{ km}$  (the *Moho* discontinuity)

2 geographical spreading of the sources is very large

3 randomness of differential movement at a given location depends strongly on the **fractured state** of the site (surface behaves as a number of independent blocks excited from below)

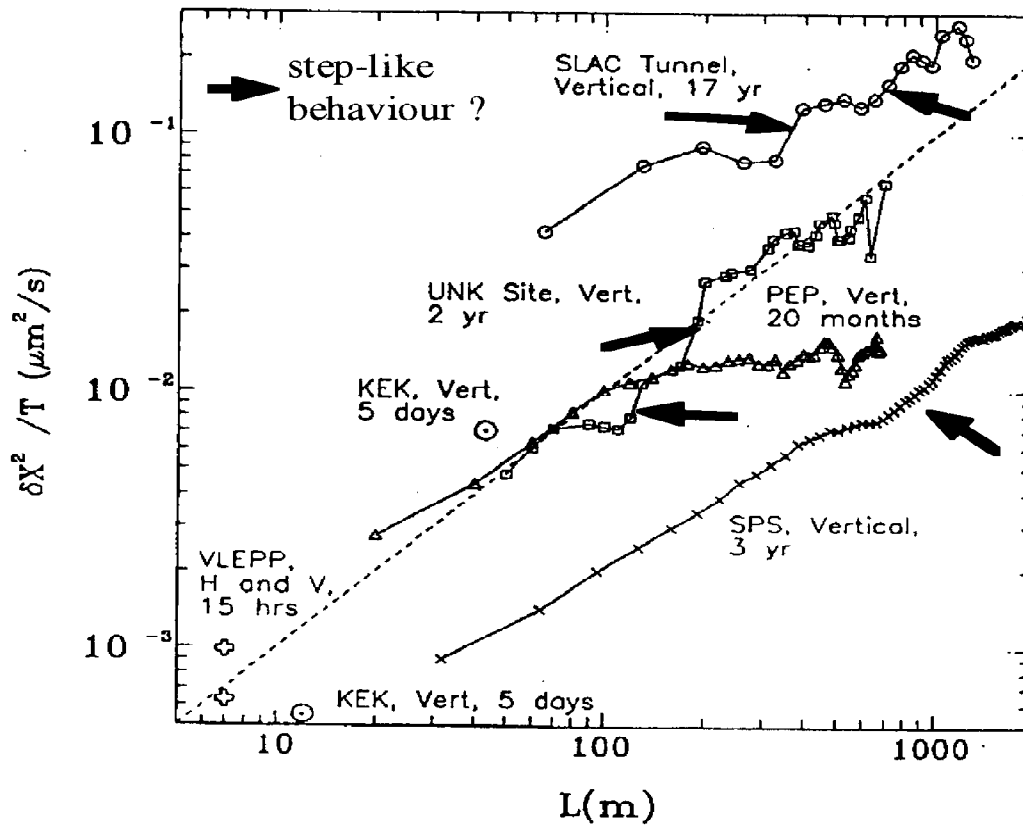
last point borne out clearly by the *experimental observation* on two points on either side of a **construction joint**.

- **introduce the statistical notion of coherence length  $L_{ch}$  :**

two points at a distance smaller than  $L_{ch}$  are **likely** to move *coherently*, while two points which are further away are **likely** to move *incoherently*

notion of coherent length is well suited for accelerators where local differences average out

indications about *coherence length* can be found in this figure (Parkhomchuk e.a. 1994) :



Differential movement as a function of distance in a number of machines for a several observation periods.

power of the differential motion seems to jump to larger values **above a given distance** ( *coherence length* ?) in each site

$L_{ch}$  seems to vary from 100...200  $\text{m}$  (UNK) to 600  $\text{m}$  (SPS)

### 3 OBSERVATIONS WITH BEAM

coherence length can be determined from orbit measurements assuming an average value for  $k_{gm} = 10^{-18} \text{ m}^2/\text{s}^2$ .

integration noise spectrum yields **displacement**<sup>2</sup> of a single element :

$$dx^2(t) = \frac{k_{gm}}{f_{co}^2} \left( \sqrt{1 + (f_{co}t)^2} - 1 \right).$$

**orbit deformation** : multiply with an **optical amplification** factor :

$$\Psi_{O_A} = (\beta Kl / 2 \sin(\pi q))^2 N$$

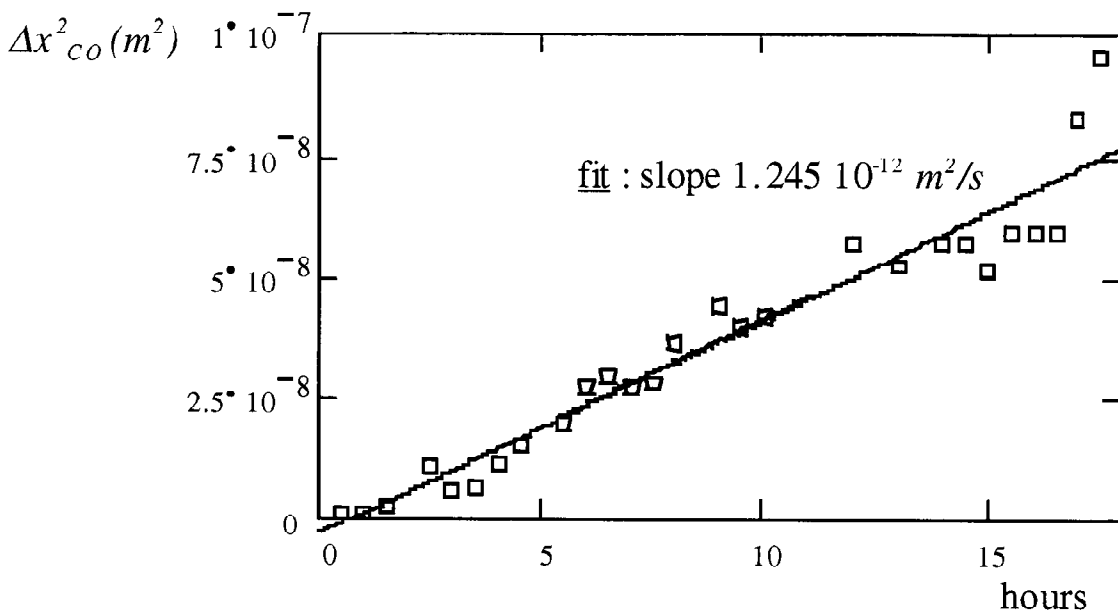
- $\beta$  the optical function at a quadrupole
- $Kl$  integrated focalisation force
- $N$  is the number of **uncorrelated** blocks around the accelerator (at the maximum the number of  $F$  or  $D$  quadrupoles)

Brinkman & Rossbach (1994)

**HERA**-proton  $L_{ch} : 250 \text{ m}$   
**HERA**-electron  $L_{ch} : 280 \text{ m}$

Tecker (1996)

**LEP**  $L_{ch} : 130 \text{ m}$   
 (known effect supra quads removed)



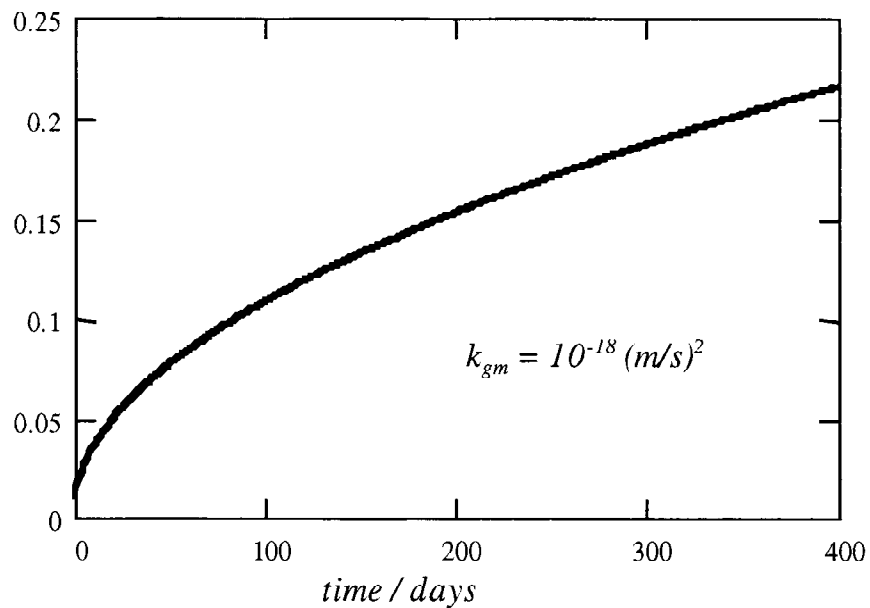
## 4 APPLICATION TO LHC

- basic equation 
$$dx^2(t) = \frac{k_{gm}}{f_{co}^2} \left( \sqrt{1 + (f_{co}t)^2} - 1 \right)$$

spans a **large time/frequency** scale

optimistic example of long term misalignment assuming quiet conditions during 1 year! :

*differential  
displacement / mm*

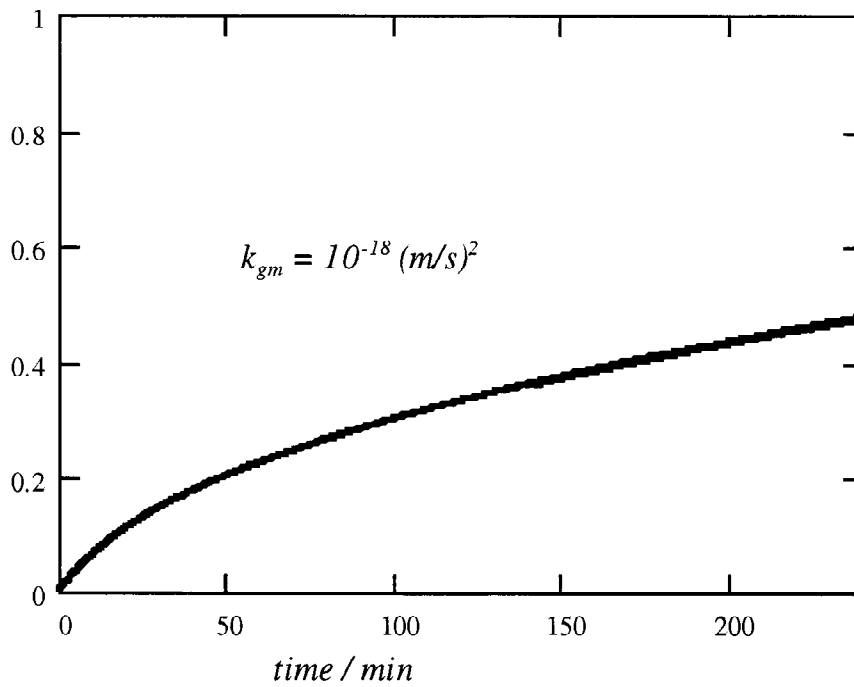


- computation of the *rms* half separation between the beams as a function of time (expressed in terms of the *rms* beam size  $\sigma$ )

**Parameters :**

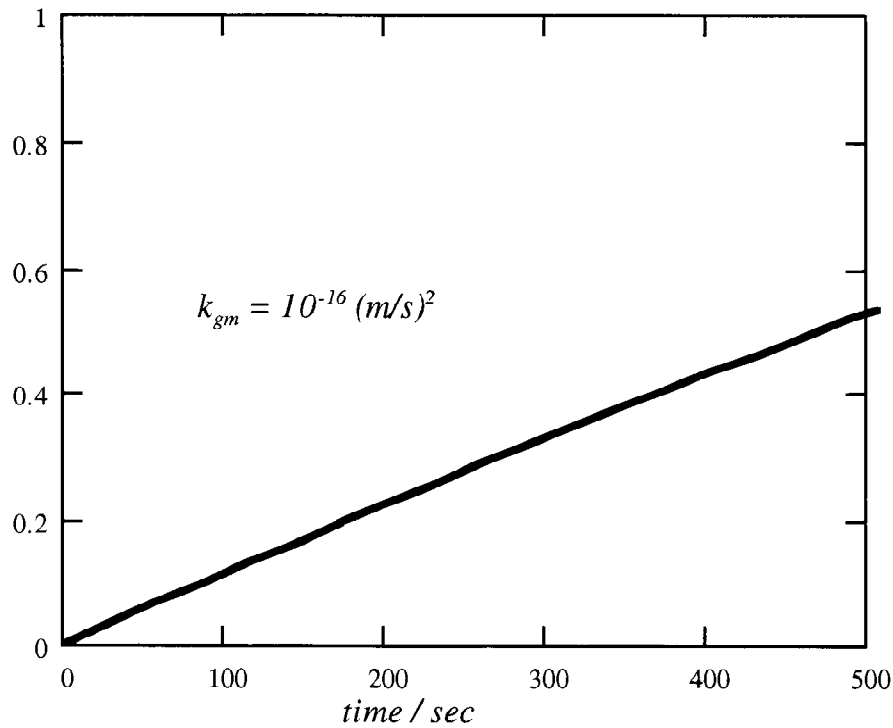
$\beta = 183$ ,  $Kl = 0.03 \text{ m}^{-1}$ ,  $q = 0.31$ ,  $N_c = 248$  (arc +DS),  $L_c = 107 \text{ m}$ ,  
 $L_{ch} = 130 \text{ m}$ , transverse emittance  $3.75 \text{ } \mu\text{radm}$ ,  $7 \text{ TeV}/c$ .

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- to be expected that once in a while the global system is *highly excited* ( $k_{gm} = 10^{-16} \text{ m}^2/\text{s}^2$ ):

half sep. /  $\sigma$



====> procedures must be ready to cope with such a rate of separation

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